

# Structural Optimization of 3D Masonry Buildings (Supplemental Material)

## A Matrix Structure

We detail the matrix equation for static equilibrium:

$$\mathbf{A}_{eq} \cdot \mathbf{f} + \mathbf{w} = \mathbf{0}$$

$\mathbf{w}_j$ :  $6 \times 1$  vector containing the 3D weight and net torque for block  $j$ . Typically the only non-zero element is the  $z$ -component of weight. For any external loads acting on block  $j$ , the force and torque contributions are added here.

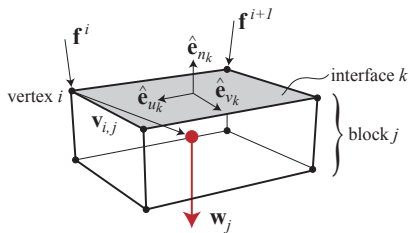
$\mathbf{r}_k$ : Contains the unknown force vectors  $\mathbf{f}^i$ , for vertices  $i$  on interface  $k$ .  $height(\mathbf{r}_k)$  is  $4v_k$ , where  $v_k$  is the number of vertices on interface  $k$  and each vertex contributes a 3D force with positive and negative parts for the axial forces.

$\mathbf{A}_{j,k}$ : Submatrices  $\mathbf{A}_{j,k}$  contain coefficients for net force and net torque contributions from interface  $k$  acting on block  $j$ . Each  $\mathbf{A}_{j,k}$  has dimension  $6 \times height(\mathbf{r}_k)$ . Rows 1-3 are coefficients for net force contributions in  $x, y, z$  and rows 4-6 are coefficients for net torque contributions about the  $x, y, z$  axes.

$$\mathbf{A}_{j,k} \mathbf{r}_k = \begin{bmatrix} \mathbf{F}_k & \mathbf{F}_k & \dots \\ \mathbf{T}_{i,j,k} & \mathbf{T}_{i+1,j,k} & \dots \end{bmatrix} \begin{bmatrix} \mathbf{f}^i \\ \vdots \end{bmatrix}$$

$\mathbf{F}_k = [\hat{\mathbf{e}}_{n_k} \hat{\mathbf{e}}_{u_k} \hat{\mathbf{e}}_{v_k}]$  and  $\mathbf{T}_{i,j,k} = [(\hat{\mathbf{e}}_{n_k} \times \mathbf{v}_{i,j}) (\hat{\mathbf{e}}_{u_k} \times \mathbf{v}_{i,j}) (\hat{\mathbf{e}}_{v_k} \times \mathbf{v}_{i,j})]$ . Unit vectors  $\hat{\mathbf{e}}_{n_k}$ ,  $\hat{\mathbf{e}}_{u_k}$  and  $\hat{\mathbf{e}}_{v_k}$  are the normal vector and friction basis vectors for face  $k$  (see Figure 1). The subscript  $x$  refers to the  $x$ -component of the vector.

The number of submatrices  $\mathbf{A}_{j,k}$  in row  $j$  of  $\mathbf{A}_{eq}$  is equal to the number of neighbors incident on block  $j$ . There are two submatrices in each column  $k$ , since  $\mathbf{r}_k$  represents the interaction between surfaces of two adjacent blocks.



**Figure 1:** Indexing for equations of static equilibrium. Vector  $\hat{\mathbf{e}}_{n_k}$  is the unit normal for interface  $k$ , and  $\hat{\mathbf{e}}_{u_k}$  and  $\hat{\mathbf{e}}_{v_k}$  are the directions of in-plane friction forces. Vector  $\mathbf{v}_{i,j}$  is the relative position of vertex  $i$  w.r.t. the centroid of block  $j$ .  $\mathbf{w}_j$  is the 3D weight vector for block  $j$ .

## B Partial Derivatives

The partial derivatives of coefficients for net force equilibrium,  $\mathbf{F}_k$ , on face  $k$  are:

$$\frac{\partial \mathbf{F}_k}{\partial \omega} = \frac{\partial [\hat{\mathbf{e}}_n \hat{\mathbf{e}}_u \hat{\mathbf{e}}_v]_k}{\partial \omega}$$

where  $\omega$  is a parameter from the set  $\{u_{i,k}, v_{i,k}, n_k, \theta_k, \phi_k\}$  as described in §6. The partial derivatives of coefficients for net torque equilibrium,  $\mathbf{T}_k$ , on face  $k$  are:

$$\begin{aligned} \frac{\partial \mathbf{T}_{i,j,k}}{\partial \omega} &= \frac{\partial [(\hat{\mathbf{e}}_n \times \mathbf{v}_{i,j}) (\hat{\mathbf{e}}_u \times \mathbf{v}_{i,j}) (\hat{\mathbf{e}}_v \times \mathbf{v}_{i,j})]_k}{\partial \omega} \\ \frac{\partial (\hat{\mathbf{e}}_n \times \mathbf{v}_{i,j})_k}{\partial \omega} &= \hat{\mathbf{e}}_n \times \frac{\partial \mathbf{v}_{i,j}}{\partial \omega} + \frac{\partial \hat{\mathbf{e}}_n}{\partial \omega} \times \mathbf{v}_{i,j} \\ &= \hat{\mathbf{e}}_n \times \left( \frac{\partial \mathbf{p}_{i,j}}{\partial \omega} - \frac{\partial \mathbf{c}_j}{\partial \omega} \right) + \frac{\partial \hat{\mathbf{e}}_n}{\partial \omega} \times \mathbf{v}_{i,j} \end{aligned}$$

The derivative of the centroid position  $\mathbf{c}_j$  for block  $j$  is:

$$\frac{\partial \mathbf{c}_j}{\partial \omega} = \frac{\partial}{\partial \omega} \left( \frac{\sum v_{T_{i,j}} \mathbf{c}_{T_{i,j}}}{\sum v_{T_{i,j}}} \right)$$

where  $v_{T_{i,j}}$  is the volume of tetrahedron  $i$  on block  $j$  and  $\mathbf{c}_{T_{i,j}}$  is the centroid of tetrahedron  $i$ .

$$\begin{aligned} \frac{\partial v_{T_{i,j}}}{\partial \omega} &= \frac{1}{6} \text{sign}(\mathbf{a}_0 \cdot (\mathbf{a}_1 \times \mathbf{a}_2)) \frac{\partial}{\partial \omega} (\mathbf{a}_0 \cdot (\mathbf{a}_1 \times \mathbf{a}_2)) \\ \frac{\partial \mathbf{c}_{T_{i,j}}}{\partial \omega} &= \frac{1}{4} \frac{\partial}{\partial \omega} (\mathbf{a}_0 + \mathbf{a}_1 + \mathbf{a}_2) \end{aligned}$$

where coordinates  $\mathbf{a}_0, \mathbf{a}_1, \mathbf{a}_2$  are three corners of tetrahedron  $i$ , offset such that the fourth coordinate lies at the origin  $(0, 0, 0)$ .

The derivative of the weight vector  $\partial \mathbf{w} / \partial \omega$  for block  $j$  is given by:

$$\frac{\partial \mathbf{w}_j}{\partial \omega} = \rho \frac{\partial v_j}{\partial \omega} \hat{\mathbf{g}} = \rho \left( \sum_i \frac{\partial v_{T_{i,j}}}{\partial \omega} \right) \hat{\mathbf{g}}$$

where  $\rho$  is the block density and  $\hat{\mathbf{g}}$  is the direction of gravity, and  $v_j$  is the volume of block  $j$ .

### B.1 Constraint Derivatives

In the closed form solution of  $\mathbf{f}^*$ , the constraint matrix  $\mathbf{C}$  is a concatenation of the matrix  $\mathbf{A}_{eq}$  (static equilibrium), the active friction-cone inequalities of the matrix  $\mathbf{A}_{fr}$ , and the active lower bounds on  $\mathbf{f}$ . The friction constraints and lower bound constraints are not dependent on block geometry, giving  $\partial \hat{\mathbf{A}}_{fr} / \partial \omega = \mathbf{0}$  and  $\partial \tilde{\mathbf{1}}_{lb} / \partial \omega = \mathbf{0}$ .

$$\frac{\partial \mathbf{C}}{\partial \omega} = \begin{bmatrix} \partial \mathbf{A}_{eq} / \partial \omega \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}, \quad \frac{\partial \mathbf{b}}{\partial \omega} = \begin{bmatrix} -\partial \mathbf{w} / \partial \omega \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}$$

where  $\omega$  is a parameterization of the structure's geometry  $\Omega$ . The derivations of  $\partial \mathbf{A}_{eq} / \partial \omega$  and  $\partial \mathbf{w} / \partial \omega$  are shown above.

### B.2 Energy Derivatives

The expression for the derivative of  $\mathbf{f}^*$  is then obtained by differentiating the closed form expression:

$$\frac{\partial \mathbf{f}^*_{\Omega}}{\partial \omega} = \mathbf{H}^{-1} \left( \frac{\partial \mathbf{C}^T}{\partial \omega} \mathbf{E}^{-1} \mathbf{b} + \mathbf{C}^T \mathbf{E}^{-1} \left( \frac{\partial \mathbf{b}}{\partial \omega} - \frac{\partial \mathbf{E}}{\partial \omega} \mathbf{E}^{-1} \mathbf{b} \right) \right)$$

$\mathbf{H}$  is the weighting matrix for the objective function and is held constant, and  $\mathbf{E} = \mathbf{C}\mathbf{H}^{-1}\mathbf{C}^T$ . The derivative of  $\mathbf{E}$  by application of the chain rule is:

$$\partial\mathbf{E}/\partial\omega = \partial\mathbf{C}/\partial\omega \mathbf{H}^{-1}\mathbf{C}^T + \mathbf{C}\mathbf{H}^{-1}\partial\mathbf{C}^T/\partial\omega$$

The terms  $\partial\mathbf{C}/\partial\omega$  and  $\partial\mathbf{b}/\partial\omega$  describe how the constraints change as the geometry changes according to parameterization  $\omega$ . The expression for the gradient of  $y(\Omega)$  is given by:

$$\nabla y = \alpha \nabla y_{uniform} + \nabla y_{torque}$$

where the derivatives of uniform and torque tension energies are:

$$\begin{aligned} \frac{\partial y_{uniform}}{\partial\omega} &= \mathbf{f}^{*T} \mathbf{H}_{uniform} \frac{\partial \mathbf{f}^*}{\partial\omega} \\ \frac{\partial y_{torque}}{\partial\omega} &= \mathbf{f}^{*T} \mathbf{H}_{torque} \frac{\partial \mathbf{f}^*}{\partial\omega} + \frac{1}{2} \left( \mathbf{f}^{*T} \frac{\partial \mathbf{H}_{torque}}{\partial\omega} \mathbf{f}^* \right) \end{aligned}$$

with  $\partial\mathbf{H}_{torque}/\partial\omega = (\mathbf{I} - \mathbf{H}_{min})^T \partial\mathbf{D}_{torque}/\partial\omega (\mathbf{I} - \mathbf{H}_{min})$ . Matrices  $\mathbf{H}_{uniform}$  and  $\mathbf{H}_{min}$  are constant since it assumed minimum-tension vertices remain the same for differential movement.

## C Cables

For gradient computation, we parametrize cables using the x,y,z coordinates of their end points. The derivative of the weight vector  $\partial\mathbf{w}/\partial\omega$  for cable is given by:

$$\frac{\partial\mathbf{w}}{\partial\omega} = \rho \frac{\partial L}{\partial\omega} \hat{\mathbf{g}}$$

where  $L$  is the length of the cable, and  $\rho$  is its the mass per unit length.

$$\frac{\partial L}{\partial\omega} = \frac{1}{2} \frac{(\mathbf{p}_0 - \mathbf{p}_1) \cdot \frac{\partial \mathbf{p}_0}{\partial\omega}}{\|\mathbf{p}_0 - \mathbf{p}_1\|}$$

where coordinates  $\mathbf{p}_0, \mathbf{p}_1$  are two ends of the cable.

The derivative of the cable tension direction is

$$\begin{aligned} \frac{\partial \hat{\mathbf{e}}_t}{\partial\omega} &= \frac{\partial}{\partial\omega} \left( \frac{\mathbf{p}_0 - \mathbf{p}_1}{\|\mathbf{p}_0 - \mathbf{p}_1\|} \right) \\ &= \frac{\|\mathbf{p}_0 - \mathbf{p}_1\| \cdot \left( \frac{\partial \mathbf{p}_0}{\partial\omega} \right) - (\mathbf{p}_0 - \mathbf{p}_1) \cdot \left( \frac{\partial \|\mathbf{p}_0 - \mathbf{p}_1\|}{\partial\omega} \right)}{\|\mathbf{p}_0 - \mathbf{p}_1\|^2} \end{aligned}$$